# Interlayer impact excitation by hot electrons in an

# atomic layer semiconductor heterostructure

Fatemeh Barati<sup>1,2,4,†</sup>, Max Grossnickle<sup>1,2,4,†</sup>, Shanshan Su<sup>1,3,4</sup>, Roger Lake<sup>1,3,4</sup>, Vivek Aji<sup>1</sup>, and Nathaniel M. Gabor<sup>1,2,4</sup>

<sup>1</sup>University of California at Riverside Department of Physics and Astronomy <sup>2</sup>Quantum Materials Optoelectronics Laboratory <sup>3</sup>Laboratory for Terascale and Terahertz Electronics (LATTE) <sup>4</sup>SHINES Energy Frontier Research Center





We fabricated a device from MoSe<sub>2</sub> and WSe<sub>2</sub> using exfoliated flakes. The device had source and drain contacts, as well as a back gate.

2L-WSe<sub>2</sub>

MoSe<sub>2</sub>

Quantum Materials Optoelectronics Lab



2L-WSe<sub>2</sub>

#### Discussion

Transition metal dichalcogenides (TMDs) exhibit optoelectronic properties that vary strongly with sample thickness. By using them in Van der Waals heterostructures, we can custom-build novel devices perfectly tailored to the study of interlayer electron-hole pair (exciton) generation. Efficient generation of interlayer excitons could be used to make novel photodetectors, electroluminescent emitters, or excitonic integrated circuits.

In this work, we observed efficient multiplication of interlayer excitons by hot electron relaxation across the interface of a Van der Waals heterostructure. Electronic transport measurements showed large negative differential conductance, source-drain and gate voltage dependent interlayer current, and strong temperature dependence, all of which support hot electron relaxation as the interlayer exciton generation mechanism.



(a.u.)

Photoluminescence of the overlap area shows significant suppression in the peaks when compared to the original flakes. The differences indicate that there is an additional electron relaxation mechanism.



We used 2L-WSe<sub>2</sub> and ML-MoSe<sub>2</sub> in order to study electron transfer at the interface of this heterostructure.



# **Temperature dependence**

The magnitude of the reverse bias interlayer current rapidly increases with temperature, with the maximum current at 360K being ten times as large as the maximum at 310K. Above 340K, NDC appears at high gate voltages. Asymmetry **7**-0.8 between forward and reverse bias at is also -1.2 especially noticeable at 360K.

Rescaling the reverse bias current as  $Tln(I/T^2)$ , collapses data at all temperatures to a single temperature-independent characteristic. Rescaling the forward bias current as Tln(I/T) similarly accounts for the temperature dependence.







### **IV characteristics**



## Impact excitation model

$$e_{W} + Ke \rightarrow eMo + (e_{Mo} + hW)$$

$$\stackrel{\varepsilon(k)}{\downarrow} \stackrel{\uparrow}{\downarrow} \stackrel{\uparrow}{\downarrow} \stackrel{\uparrow}{\downarrow} \stackrel{\downarrow}{\downarrow} \stackrel{\uparrow}{\downarrow} \stackrel{\downarrow}{\downarrow} \stackrel{\downarrow}{$$

The turn-on point of the device is identified by the maximum in  $dlog l/dV_{G}$ .

In forward bias, the turn-on occurs at a constant Vg S but tilts in reverse bias.

This behavior is almost

VsD

### **Device turn-on**



V<sub>SD</sub>\* labels the intercept of the linear fits of the maxima of dlogl/dVg and V<sup>+</sup><sub>ON</sub>.

V<sup>\*</sup><sub>SD</sub>→

V<sup>\*</sup><sub>SD</sub>→

At voltages above  $V_{SD}^*$ , interlayer transport occurs as low-energy electrons from the WSe<sub>2</sub> conduction



